

Normalizations on CHEASE

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1 CHEASE normalizations, MKSA units

From

$$\underline{B} = T \underline{\nabla} \phi + \underline{\nabla} \phi \times \underline{\nabla} \psi \quad (1)$$

then

$$[B] = \frac{[T]}{[R]} \rightarrow T_{\text{CHEASE}} = \frac{T_{\text{phys}}}{\overline{B}_0 \overline{R}_0} \quad (2)$$

Where R_0 and B_0 are the values used for normalizing the parameters

$$B_{0\text{CHEASE}} = B_{0\text{phys}}/B_0 \quad R_{0\text{CHEASE}} = R_{0\text{phys}}/R_0 \quad (3)$$

If one uses the boundary conditions $T_{\text{CHEASE}}(\text{edge}) = 1$, then B_0 is the vacuum magnetic field at $R = R_0$ such that $T(\text{edge}) = R_0 B_0$. From

$$\underline{\nabla} \times \underline{B} = \mu_0 \underline{j} \quad (4)$$

we obtain the following expression

$$[j] = \frac{[B]}{\mu_0 [R]} = \frac{[B]}{\mu_0 [R]} \quad (5)$$

which gives the following normalization for the current density:

$$j_{\text{CHEASE}} = j_{\text{phys}} \frac{\mu_0 R_0}{B_0} \quad (6)$$

and the corresponding relation for total current

$$[I] = \frac{[B] [R]}{\mu_0} \rightarrow I_{\text{CHEASE}} = I_{\text{phys}} \frac{\mu_0}{R_0 B_0} \quad (7)$$

For the pressure, from

$$\nabla p = j \times B \quad (8)$$

we obtain

$$[p] = [j] [B] [R] = \frac{B_0}{\mu_0 R_0} B_0 R_0 = \frac{B_0^2}{\mu_0} \quad (9)$$

Then we define

$$p_{\text{CHEASE}} = p_{\text{phys}} \frac{\mu_0}{B_0^2} \quad (10)$$

The magnetic field itself is defined by Eq. 1 which brings to

$$[\psi] = [B] [R]^2 \quad (11)$$

It follows

$$\psi_{\text{CHEASE}} = \frac{\psi_{\text{phys}}}{B_0 R_0^2} \quad (12)$$

In the same way

$$p' = \left[T \frac{\partial p}{\partial \psi} \right] = \frac{[p]}{[\psi]} = \frac{B_0^2}{\mu_0 R_0^2 B_0} = \frac{B_0}{\mu_0 R_0^2} \rightarrow p'_{\text{CHEASE}} = p'_{\text{phys}} \frac{\mu_0 R_0^2}{B_0} \quad (13)$$

and

$$TT' = \left[T \frac{\partial T}{\partial \psi} \right] = \frac{[T]^2}{[\psi]} = \frac{B_0^2 R_0^2}{B_0 R_0^2} = B_0 \rightarrow TT'_{\text{CHEASE}} = \frac{TT'_{\text{phys}}}{B_0} \quad (14)$$

Let us now check β_{pol} as given in Table 1 of H. Lutjens et al, Comput. Phys. Commun. 97 (1996) 219:

$$\beta_{\text{pol}} = \left[\frac{8\pi}{I_\phi^2 R_0} \bar{p} V_{\text{tot}} \right]_{\text{CHEASE}} \quad (15)$$

Note that $V_{\text{tot,CHEASE}}$ is in fact $V_{\text{CHEASE}}/(2\pi)$, thus we get after the above transformations:

$$\beta_{\text{pol}} = \frac{4}{\mu_0} \frac{\int p_{\text{phys}} dV_{\text{phys}}}{I_{\phi_{\text{phys}}}^2 R_{0_{\text{phys}}}} \quad (16)$$

which corresponds to the usual definition with $\overline{B}_{\text{pol}} = \mu_0 I_{\phi_{\text{phys}}} / (2V/R_{0,\text{phys}})^{1/2}$

Let us now focus on the parallel component of the plasma current as given in Eq. (8) of the above paper. We get:

$$[I_{\parallel}] = \frac{[j]}{[1/R]} = [j] [R] \quad (17)$$

Instead of $[j]$. This is because a $1/R_0$ is missing in Eq. (8). R_0 is supposed to be equal to 1 in CHEASE, so it's not so important, however the correct definition should read and has been modified in CHEASE:

$$I_{\parallel} = \frac{\langle \mathbf{j} \cdot \mathbf{B} \rangle}{R_0 \langle \mathbf{B} \cdot \nabla \phi \rangle} \quad (18)$$

In this way, I_{\parallel} has the correct dimension and we have:

$$I_{\parallel} = [j] \Rightarrow I_{\parallel \text{CHEASE}} = I_{\parallel \text{phys}} \frac{\mu_0 R_0}{B_0} \quad (19)$$

2 CHEASE normalizations in CGS

In CGS

$$\underline{\nabla} \times B = \frac{4\pi j}{c} + \frac{1}{c} \frac{\partial E}{\partial t} \quad (20)$$

$$mn \frac{dv}{dt} = qnE + \frac{qnv}{c} \times B - \underline{\nabla} p \quad (21)$$

$$\Rightarrow \frac{j}{c} \times B = \underline{\nabla} p \quad (22)$$

$$\underline{B} = T \underline{\nabla} \phi + \underline{\nabla} \phi \times \underline{\nabla} \psi \quad (23)$$

Dimensionally, again

$$[T] = [B_0] [R_0] \quad (24)$$

where this time B_0 [Gauss] and R_0 [cm]. The latter relations thus bring to

$$[j]_{\text{CGS}} = \frac{c}{4\pi} \frac{[B]}{[R]} = \frac{c}{4\pi} \frac{[B_0]}{[R_0]} \quad (25)$$

$$[p]_{\text{CGS}} = \frac{[j] [B]}{c} [R] = \frac{c}{4\pi} \frac{B_0}{R_0} \frac{1}{c} B_0 R_0 = \frac{B_0^2}{4\pi} \quad (26)$$

$$[\psi] = [B] [R]^2 = B_0 R_0^2 \quad (27)$$

$$\Rightarrow [p']_{\text{CGS}} = \frac{B_0^2}{4\pi B_0 R_0^2} = \frac{B_0}{4\pi R_0^2}$$

$$\Rightarrow p'_{\text{CHEASE}} = p'_{\text{CGS}} \frac{4\pi R_{0\text{CGS}}^2}{B_{0\text{CGS}}} \quad (28)$$

Finally we can express TT' from these relations as

$$TT'_{\text{CHEASE}} = \frac{TT'_{\text{CGS}}}{B_{0\text{CGS}}} \quad (29)$$

Note that the following identity holds

$$\begin{aligned} [p]_{\text{MKS}} &= \frac{B_{0\text{MKS}}^2}{\mu_0} = \frac{B_{0\text{MKS}}^2}{4\pi} 10^7 \\ [p]_{\text{CGS}} &= \frac{B_{0\text{CGS}}^2}{4\pi} = \frac{B_{0\text{MKS}}^2}{4\pi} 10^8 \Rightarrow [p]_{\text{CGS}} = 10 [p]_{\text{MKS}} \end{aligned} \quad (30)$$

which is the conversion factor in the NRL plasma formulary.